

Introduction to Control (00340040)

lecture no. 2

Leonid Mirkin

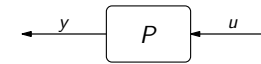
Faculty of Mechanical Engineering
Technion—IIT



1/36

Abstract control problem to begin with

Setup:



where

- P is a plant
may comprise actual controlled process, actuators, sensors, et cetera
- u is a control signal (control input)
- y is a controlled (regulated) signal (output)

Problem: Given P , find u resulting in a desired y .

2/36

Outline

Modeling

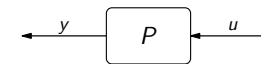
Case-study systems and their first-principles modeling

Actuation: DC motors and their first-principles modeling

Parameter identification of DC motor

3/36

What is it about?



Model is a

- description of systems using an abstract (e.g. mathematical) language.

Modeling lets us handle problems of various nature, e.g.

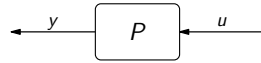
- mechanical,
- electrical,
- biological,
- social,
- ...

in a unified manner. It must be realized though that

- **models** of real-world phenomena are **never perfect**,
they are just (more or less accurate) approximations of real processes.

4/36

Why to model?



From control viewpoint, modeling is necessary just because

- if we do not know how the plant responds to our actions, then control tasks are hopeless.

In other words,

- **model-free control** is essentially a **coin tossing**.

How to model?

Essentially, three ways:

1. from first principles *ab initio*
2. phenomenological e.g. predator-prey, SIR
3. from observing experimental I/O relations *identification*

5/36

Outline

Modeling

Case-study systems and their first-principles modeling

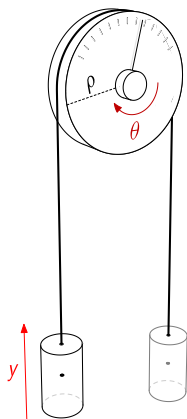
Actuation: DC motors and their first-principles modeling

Parameter identification of DC motor

6/36

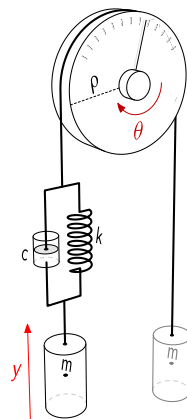
Case-studies

System Σ_1 :



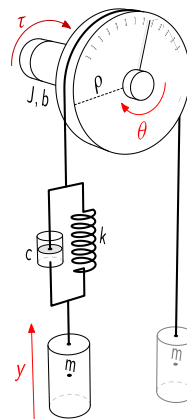
- control mass position y
- via pulley angle θ
- elevator as motivation

System Σ_2 :



- control mass position y
- via pulley angle θ
- elevator with long hoistway as motivation

System Σ_3 :

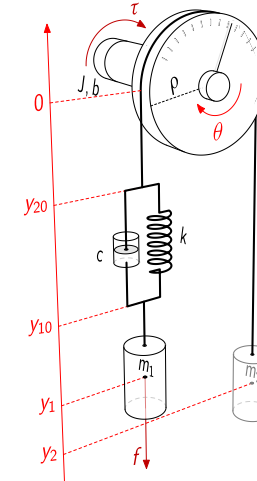


- control mass position y
- via torque τ

7/36

Data

Consider



where

- ρ : pulley radius
- m_1, m_2 : masses
- k : spring constant
- c : damping coefficient
- J : moment of inertia of pulley + shaft
- b : friction coefficient of pulley + shaft
- f : external force applied to m_1 (2nd input)

Denote also

- f_{t1}, f_{t2} : tension forces at m_1 and m_2 , respectively
- $L_0 := y_{20} - y_2$ at $\theta = 0$
- $L_1 := y_{10} - y_1$
- $L_2 := \pi\rho - y_{20} - y_2$
- $L := L_1 + L_2$ length of non-elastic part

8/36

Motion equations

Masses dynamics:

$$-m_1\ddot{y}_1(t) = m_1g + f(t) - f_{t1}(t) \quad \text{and} \quad -m_2\ddot{y}_2(t) = m_2g - f_{t2}(t)$$

Spring-damper dynamics:

$$c(\dot{y}_{20}(t) - \dot{y}_{10}(t)) + k(y_{20}(t) - y_{10}(t)) = f_{t1}(t)$$

Pulley dynamics:

$$J\ddot{\theta}(t) + b\dot{\theta}(t) = \tau(t) - \rho f_{t1}(t) + \rho f_{t2}(t)$$

Algebraic constraints:

$$\begin{aligned} y_{10}(t) - y_1(t) &= L_1, \\ -y_{20}(t) - y_2(t) &= L_2 - \pi\rho, \\ y_{20}(t) - y_2(t) &= L_0 + 2\rho\theta(t) \end{aligned}$$

9/36

Elimination of variables

From algebraic constraints,

$$y_{10} = y_1 + L_1, \quad y_{20} = \rho\theta + \frac{L_0 - L_2 + \pi\rho}{2}, \quad y_2 = -\rho\theta - \frac{L_0 + L_2 - \pi\rho}{2}$$

Then, spring-damper verifies

$$c(\rho\dot{\theta} - \dot{y}_1) + k(\rho\theta - y_1) + k\left(\frac{L_0 - L_2 + \pi\rho}{2} - L_1\right) = f_{t1}$$

and the tension forces are

$$\begin{aligned} f_{t1} &= m_1\ddot{y}_1 + m_1g + f, \\ f_{t2} &= m_2\ddot{y}_2 + m_2g = -\rho m_2\ddot{\theta} + m_2g \end{aligned}$$

10/36

System equations

Thus, we have:

$$\begin{cases} m_1\ddot{y}_1 + c\dot{y}_1 + ky_1 + k\left(\frac{m_1g}{k} + \frac{L + L_1 - L_0 - \pi\rho}{2}\right) = \rho(c\dot{\theta} + k\theta) - f \\ (J + \rho^2 m_2)\ddot{\theta} + b\dot{\theta} = \tau - \rho m_1\ddot{y}_1 - \rho(m_1 - m_2)g - \rho f \end{cases}$$

This is a **nonlinear** set of equations (superposition principle doesn't hold¹).

¹Just think of the simpler system $y = u + 1$. Since $(u_1 + u_2) + 1 \neq (u_1 + 1) + (u_2 + 1)$, it's nonlinear. Such systems are called *affine* and can be linearized precisely via introducing deviation variables. For example, define $\tilde{u} = u + 1$, which yields linear system $y = \tilde{u}$.

11/36

Equilibrium

In equilibrium, the system satisfies algebraic equations

$$\begin{cases} y_1 = \rho\theta - \frac{L + L_1 - L_0 - \pi\rho}{2} - \frac{f + m_1g}{k} \\ \tau = \rho(m_1 - m_2)g + \rho f \end{cases}$$

If we *assume* that $y_1 = y_2$ at $\theta = 0$ and $f = 0$, an additional constraint

$$-\frac{L + L_1 - L_0 - \pi\rho}{2} - \frac{m_1g}{k} = -\frac{L_0 + L_2 - \pi\rho}{2}$$

yields $L_0 = L_1 + m_1g/k$, so that

$$\begin{cases} y_1 = \rho\theta - \frac{L - \pi\rho + m_1g/k}{2} - \frac{f}{k} \\ \tau = \rho(m_1 - m_2)g + \rho f \end{cases}$$

We then **choose** the equilibrium corresponding to $\theta = 0$ and $f = 0$.

12/36

Linearization

Defining **deviation variables**

$$\tilde{y}_1(t) := y_1(t) + \frac{L - \pi\rho + m_1g/k}{2} \quad \text{and} \quad \tilde{\tau}(t) := \tau(t) - \rho(m_1 - m_2)g$$

we end up with the following **linear** model:

$$\begin{cases} m_1\ddot{\tilde{y}}_1(t) + c\dot{\tilde{y}}_1(t) + k\tilde{y}_1(t) = \rho(c\dot{\theta}(t) + k\theta(t)) - f(t) \\ (J + \rho^2 m_2)\ddot{\theta}(t) + b\dot{\theta}(t) = \tilde{\tau}(t) - \rho m_1\ddot{\tilde{y}}_1(t) - \rho f(t) \end{cases}$$

or

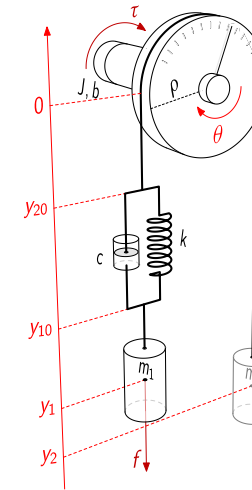
$$\begin{cases} \tilde{Y}_1(s) = \frac{1}{m_1s^2 + cs + k} ((cs + k)\rho\Theta(s) - F(s)) \\ \Theta(s) = \frac{1}{(J + \rho^2 m_2)s^2 + bs} (\tilde{T}(s) - \rho m_1s^2\tilde{Y}_1(s) - \rho F(s)) \end{cases}$$

in the Laplace domain.

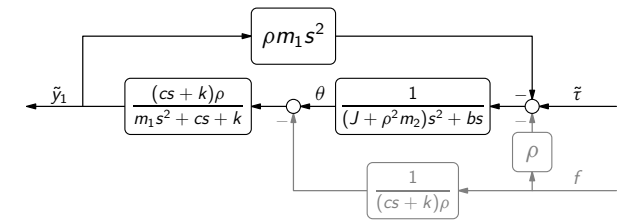
13/36

Block-diagram

Thus,



can be modeled as



in terms of deviations

$$\tilde{y}_1(t) = y_1(t) + \frac{L - \pi\rho + m_1g/k}{2},$$

$$\tilde{\tau}(t) = \tau(t) - \rho(m_1 - m_2)g$$

from the equilibrium corresponding to $y_1 = y_2$ at $\theta = 0$ and $f = 0$.

14/36

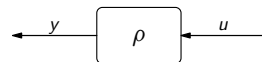
System Σ_1 (Atwood machine)

Let

Assuming inelastic string, no slippage, and $y = 0$ at $\theta = 0$,

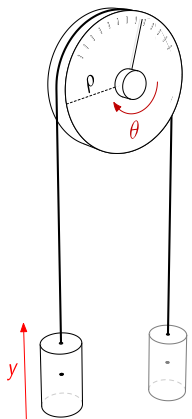
$$y(t) = \rho\theta(t) \quad (\text{with } \theta \text{ in rad}),$$

so the system can be presented by the following block diagram:



and the plant transfer function

$$P(s) = \rho$$



15/36

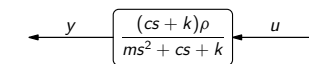
System Σ_2

Let

Now, the system motion satisfies

$$m\ddot{y}(t) + c(\dot{y}(t) - \rho\dot{\theta}(t)) + k(y(t) - \rho\theta(t)) = 0$$

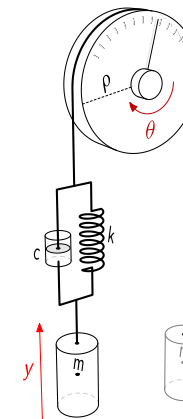
(assuming no slippage and $y = 0$ at $\theta = 0$), which leads to



and the plant transfer function

$$P(s) = \frac{(cs + k)\rho}{ms^2 + cs + k}$$

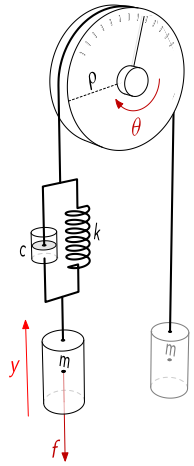
(note that $\lim_{k \rightarrow \infty} P(s) = \rho$).



16/36

System Σ_2 with disturbance

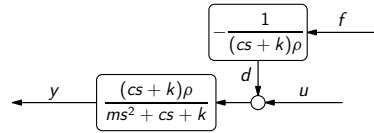
Let



Now, the system motion satisfies

$$m\ddot{y}(t) + c(\dot{y}(t) - \rho\dot{\theta}(t)) + k(y(t) - \rho\theta(t)) = -f(t)$$

(assuming no slippage and $y = 0$ at $\theta = 0$ and $f = 0$), which leads to

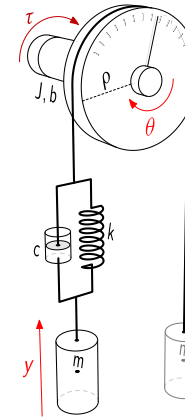


where signal d called **input disturbance**.

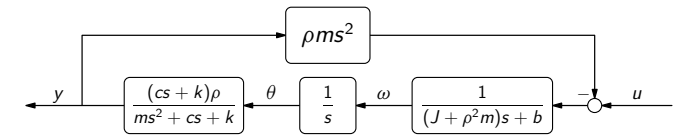
17/36

System Σ_3

Let



In this case, assuming no slippage and $y = 0$ at $\theta = 0$ we have the following block diagram:



and the plant transfer function

$$P(s) = \frac{(cs+k)\rho}{s((J+\rho^2m)s+b)(ms^2+cs+k)+\rho^2ms^2(cs+k)}$$

(it is **unstable** because of the pole at the origin).

18/36

Outline

Modeling

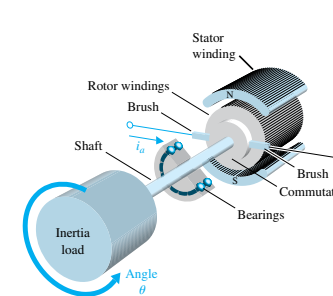
Case-study systems and their first-principles modeling

Actuation: DC motors and their first-principles modeling

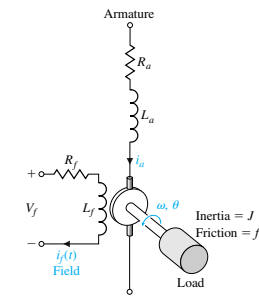
Parameter identification of DC motor

19/36

DC motors



sketch



wiring diagram

Electric motors are devices converting

- electrical energy into mechanical energy.

DC motors run on DC electric power. There are many types of DC motors, we study armature-controlled **brushed DC motors**.

20/36

DC motors (contd)

Advantages:

- high torque
- position / speed / torque controllability over a wide range
- portability
- well-behaved speed-torque characteristics
- ...

Applications (**actuators**):

- robotic manipulators
- tape transport mechanisms
- disk drivers
- ...

21/36

Modeling voltage-controlled DC motors

Important things:

1. Torque τ_m generated by the motor proportional to armature current i_a :

$$\tau_m(t) = K_m i_a(t) \quad \text{or} \quad T_m(s) = K_m I_a(s),$$

where $K_m \left[\frac{\text{Nm}}{\text{A}} \right]$ is the **motor constant** (torque constant).

2. Armature current satisfies

$$L_a \dot{i}_a(t) + R_a i_a(t) = v_a(t) - v_b(t) \quad \text{or} \quad I_a(s) = \frac{V_a(s) - V_b(s)}{L_a s + R_a},$$

where v_a is the applied input voltage and v_b is the back electromotive force (**back emf**) voltage proportional to motor angular velocity ω_m :

$$v_b(t) = K_b \omega_m(t) \quad \text{or} \quad V_b(s) = K_b \Omega_m(s),$$

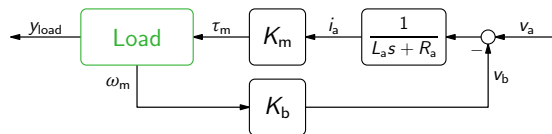
where $K_b \left[\frac{\text{Vsec}}{\text{rad}} \right]$ is the motor **back emf constant**² ($9.55 \frac{\text{Vsec}}{\text{rad}} \approx 1 \frac{\text{V}}{\text{rpm}}$).

²Normally, $K_b = K_m$ if measured in compatible units.

22/36

Modeling voltage-controlled DC motors (contd)

Resulting system can be presented as the following block diagram:



Here y_{load} is a (controlled) load output, not necessarily coinciding with the motor shaft angular velocity ω_m .

The dependence of ω_m on the load and internal feedback loop³ (back emf) renders voltage-controlled motors

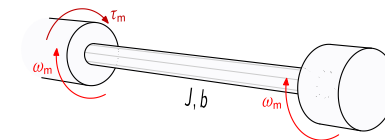
- strongly dependent on load dynamics.

We shall explicitly have ω_m as an output of the load model to incorporate the load into the motor model.

³Models substantially simplified in the *current-controlled* case (no back emf loop).

23/36

Example 1: rigid mechanical load



Consider a rigid load (e.g. the rotor itself) with $y_{\text{load}} = \omega_m$ and satisfying

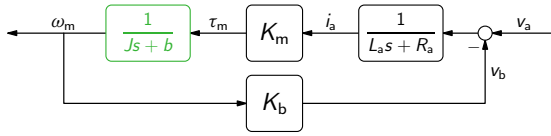
$$J \dot{\omega}_m(t) + b \omega_m(t) = \tau_m(t) \quad \text{or} \quad \Omega_m(s) = \frac{1}{Js + b} T_m(s),$$

where J is its moment of inertia and b is the friction coefficient.

24/36

Example 1: rigid mechanical load (contd)

In this case we have:



which results in the following transfer function from v_a to ω_m :

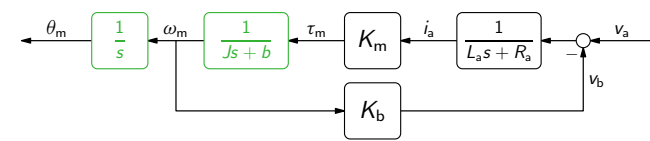
$$P_\omega(s) = \frac{K_m}{(L_a s + R_a)(J s + b) + K_b K_m},$$

which is always stable (2nd order denominator with positive coefficients).

25/36

Example 1: rigid mechanical load (contd)

If $y_{\text{load}} = \theta_m$ (motor shaft angle), then system becomes



with the transfer function

$$P_\theta(s) = \frac{K_m}{s((L_a s + R_a)(J s + b) + K_b K_m)} = \frac{1}{s} P_\omega(s),$$

which is **unstable** (pole at the origin). If armature (electrical) time constant is significantly smaller than mechanical time constant, i.e. if $\frac{L_a}{R_a} \ll \frac{J}{b}$, then

$$P_\omega(s) \approx \frac{K_m}{R_a(J s + b) + K_b K_m} \quad \text{and} \quad P_\theta(s) \approx \frac{K_m}{s(R_a(J s + b) + K_b K_m)}$$

(where L_a neglected) are sufficiently accurate.

26/36

Example 1: electrical time constant

Consider a motor (in fact, MINIMOTOR 2342) with

K_m [$\frac{\text{Nm}}{\text{A}}$]	K_b [$\frac{\text{Vsec}}{\text{rad}}$]	J [kg m^2]	b [$\frac{\text{kg m}^2}{\text{sec}}$]	R_a [Ω]	L_a [H]
0.0261	0.0261	$5.8 \cdot 10^{-7}$	$9.67 \cdot 10^{-5}$	7.1	$2.65 \cdot 10^{-4}$

with mechanical and electrical time constants of $6 \cdot 10^{-3}$ and $3.73 \cdot 10^{-5}$, respectively. This results in

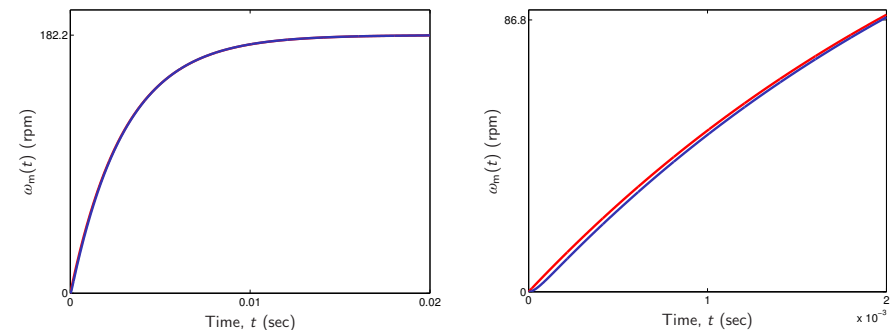
$$P_\omega(s) = \frac{19.082}{(0.003s + 1)(3.756 \cdot 10^{-5}s + 1)}$$

or, if we neglect L_a ,

$$P_\omega(s) \approx \frac{19.082}{0.003s + 1}$$

27/36

Example 1: electrical time constant (contd)



Step responses of the **second-** and **first-**order systems

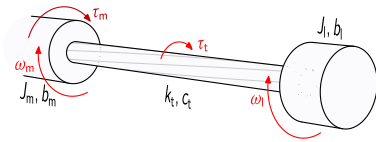
- almost indistinguishable⁴,

which justifies neglecting the dynamics of the armature circuit in this case.

⁴Except for a small difference at the start, see the close-up on the right.

28/36

Example 2: load with flexible transmission



Consider now a load (J_l and b_l) connected to the motor shaft (J_m and b_m) by a flexible inertialess transmission with the dynamics

$$\tau_t(t) = k_t \theta_\delta(t) + c_t \dot{\theta}_\delta(t) \quad \text{or} \quad T_t(s) = \frac{c_t s + k_t}{s} \Omega_\delta(s),$$

where $\theta_\delta := \theta_m - \theta_l$, $\dot{\theta}_\delta := \dot{\theta}_\delta = \omega_m - \omega_l$ and k_t and c_t are the stiffness and damping coefficient, respectively, of the transmission. The other equations:

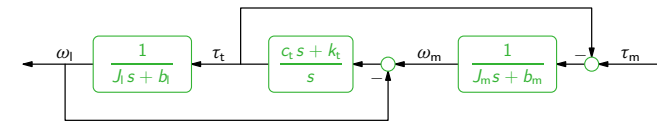
$$J_m \dot{\omega}_m(t) + b_m \omega_m(t) = \tau_m(t) - \tau_t(t) \quad \text{or} \quad \Omega_m(s) = \frac{T_m(s) - T_t(s)}{J_m s + b_m},$$

$$J_l \dot{\omega}_l(t) + b_l \omega_l(t) = \tau_t(t) \quad \text{or} \quad \Omega_l(s) = \frac{1}{J_l s + b_l} T_t(s).$$

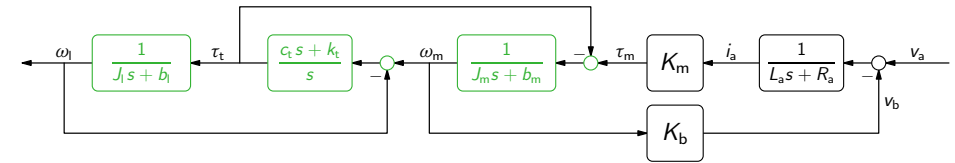
29/36

Example 2: load with flexible transmission (contd)

This system corresponds to the following block-diagram:



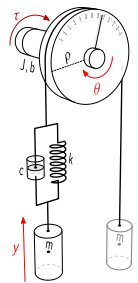
Combining load and motor block diagrams, we end up with



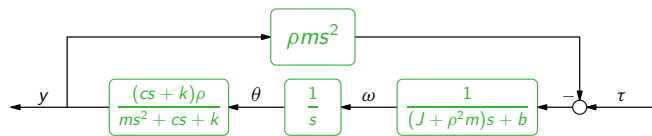
The transfer function of the system can be derived by routine block-diagram manipulations, as shown in Lecture 1.

30/36

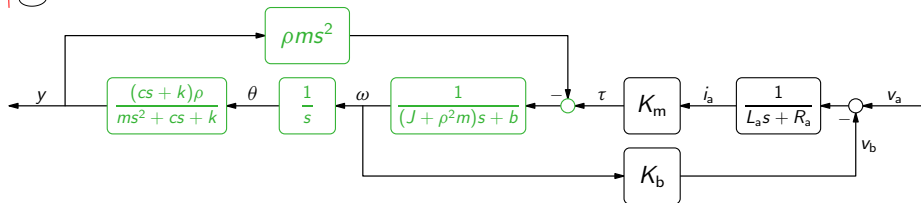
System Σ_3 as load



Remember, this system described as



If T is generated by a DC motor, we end up with the plant



having the following t.f. (assuming $L_a = 0$ and denoting $\varkappa := K_m K_b / R_a$):

$$P(s) = \frac{(cs+k)\rho K_m / R_a}{m(J+m\rho^2)s^4 + (\varkappa m + cJ + bm + 2c\rho^2)s^3 + (\varkappa c + cb + Jk + 2k\rho^2)s^2 + k(\varkappa + b)s}$$

31/36

Outline

Modeling

Case-study systems and their first-principles modeling

Actuation: DC motors and their first-principles modeling

Parameter identification of DC motor

32/36

The problem

We know that

$$P_\omega(s) \approx \frac{K_m}{R_a(Js + b) + K_b K_m}$$

but we might not know the parameters,

- some of them (K_m , K_b , R_a) can be taken from the catalog
- the others (load's J and b) are harder to calculate

Alternative to the first-principles approach:

- determining parameters from experiments (system identification)

To that end, rewrite

$$\frac{K_m}{R_a(Js + b) + K_b K_m} = \frac{k_{st}}{\tau s + 1}$$

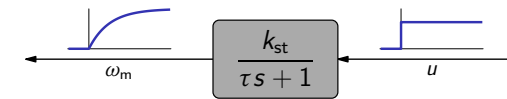
where

$$k_{st} := \frac{K_m}{K_b K_m + R_a b} \quad \text{and} \quad \tau := \frac{R_a J}{K_b K_m + R_a b}$$

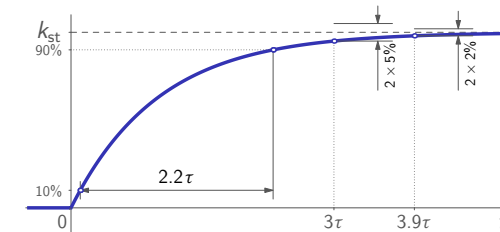
33/36

Experimental setup

We try to identify parameters k_{st} and τ from the step response



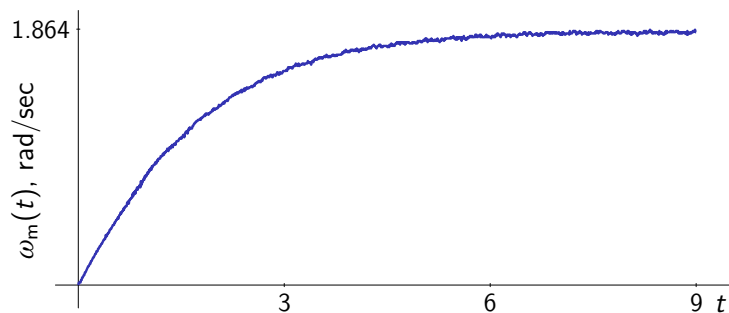
taking into account that it is relatively simple, viz. $y(t) = k_{st}(1 - e^{-t/\tau})$:



34/36

Experimental data

Reality (response to the step voltage of a magnitude of 1.2V):



is not exactly according to the theory. Reasons:

- measurement noise (sensor is an encoder, hence quantization)
- nonlinearities (e.g. mechanical friction)
- additional dynamics (inductance, eccentricity, et cetera)

Still, it closely resembles the step response of a 1-order system.

35/36

Fitting 1-order response to experimental data

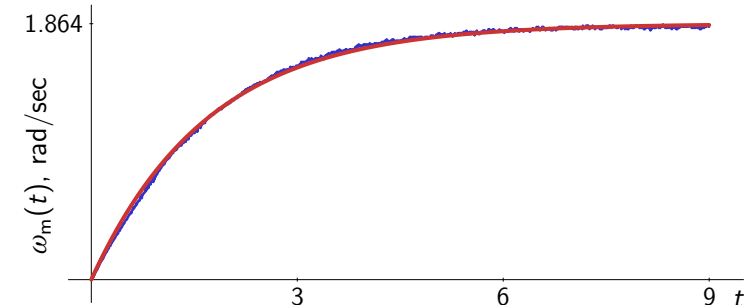
Brute-force parametric search over possible values of k_{st} and τ to fit

$$\omega_m(t) = 1.2k_{st}(1 - e^{-t/\tau})$$

into experimental data yields

$$k_{st} = 1.5533 \quad \text{and} \quad \tau = 1.7$$

with a reasonably good fit:



36/36